Recent developments in Boundary Quantum Field Theories

Z. Bajnok

Institute for Theoretical Physics, Eötvös University, Budapest

in collaboration with Changrim Ahn, G. Böhm, A. George, R. Nepomechie, L. Palla, Chaiho Rim, L. Samaj, G. Takács, Alyosha Zamolodchikov

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Relativistic Boundary QFT in 1+D (not only integrable)



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Relativistic Boundary QFT in 1+D (easy presentation) 1+1





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Relativistic Bulk QFT in 1+1



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Relativistic Bulk QFT $(x, it) \leftrightarrow (-iT, X)$



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Boundary introduced in space



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Boundary introduced in space

in time





Colloidal solution: charged or neutral particles in liquid or in gas



Colloidal solution: charged or neutral particles in liquid or in gas

Problem: Calculate the effective interaction (Hendrik Casimir 1948)



Colloidal solution: charged or neutral particles in liquid or in gas

Problem: Calculate the effective interaction (Hendrik Casimir 1948)



 $F \propto a^{-4}$



 $F \propto a^{-4}$ $\phi \propto \langle \Phi(x) \rangle = Q(\varphi_0) e^{-mx}$; $Q(\infty) < \infty$









Reduction formula







Reduction formula



S bootstrap FF bootstrap





Reduction formula



S bootstrap FF bootstrap

Reflection





Reduction formula



S bootstrap FF bootstrap

Reflection ${}^{out}\langle \theta'|\theta\rangle^{in}$ ${}^{-\theta}$

BGreen function $G_B^2(p_i, i\omega_i)$

 $(p_1,i\omega_1)$





Reduction formula



S bootstrap FF bootstrap

Reflection ${}^{out}\langle \theta'|\theta\rangle^{in}$ ${}^{-\theta}$

BGreen function $G_B^2(p_i, i\omega_i)$ $(p_1, i\omega_i)$ BReduction formula







Reduction formula



S bootstrap FF bootstrap

BReduction formula



R bootstrap BFF bootstrap



Reflection ${}^{out}\langle \theta' | \theta \rangle^{in}$

BGreen function $G_B^2(p_i, i\omega_i)$ (p_i i \omega_i) (p_i i \omega_i)





Reduction formula



S bootstrap FF bootstrap

Reflection ${}^{out}\langle \theta'|\theta\rangle^{in}$ ${}^{-\theta}$

BGreen function $G_B^2(p_i, i\omega_i)$ (p_iw_j) (p_iw_i) BReduction formula



R bootstrap BFF bootstrap









Reduction formula



S bootstrap FF bootstrap

Reflection ${}^{out}\langle \theta'|\theta\rangle^{in}$ $-\theta$

BGreen function $G_B^2(p_i, i\omega_i)$ (p_i(w_i)) (p_1,iw_i) BReduction formula



R bootstrap BFF bootstrap











Reduction formula



S bootstrap FF bootstrap

Reflection ${}^{out}\langle\theta'|\theta\rangle^{in}$ ${}^{-\theta}$

BGreen function $G_B^2(p_i, i\omega_i)$ (p_i i \omega_i) (p_i i \omega_i) BReduction formula



R bootstrap BFF bootstrap

BState $\langle B|\theta_1,\theta_2\rangle$



BReduction formula







Reduction formula



S bootstrap FF bootstrap

Reflection ${}^{out}\langle \theta' | \theta \rangle^{in}$ - θ

BGreen function $G_B^2(p_i, i\omega_i)$ $(p_2i\omega_2)$ $(p_1,i\omega_1)$ BReduction formula



R bootstrap BFF bootstrap



BState $\langle B|\theta_1,\theta_2\rangle$



BReduction formula





Bulk Hilbert space

Bulk initial state

(energy $i\omega = im \cosh \theta$, momentum $p = m \sinh \theta$)

 $|\theta_1, \theta_2, \dots, \theta_N\rangle^{in}$ $\theta_1 > \theta_2 > \dots > \theta_N$



Bulk Hilbert space

Bulk final state

 $\begin{array}{c} |\boldsymbol{\theta}_{1}^{'},\boldsymbol{\theta}_{2}^{'},\ldots,\boldsymbol{\theta}_{M}^{'}\rangle^{out} \\ \boldsymbol{\theta}_{1}^{'} < \boldsymbol{\theta}_{2}^{'} < \ldots < \boldsymbol{\theta}_{M}^{'} \end{array}$



Bulk initial state

(energy $i\omega = im \cosh \theta$, momentum $p = m \sinh \theta$)

$$|\theta_1, \theta_2, \dots, \theta_N\rangle^{in}$$

 $\theta_1 > \theta_2 > \dots > \theta_N$



Scattering matrix

Bulk final state

$$\begin{array}{c} |\boldsymbol{\theta}_{1}^{'},\boldsymbol{\theta}_{2}^{'},\ldots,\boldsymbol{\theta}_{M}^{'}\rangle^{out} \\ \boldsymbol{\theta}_{1}^{'} < \boldsymbol{\theta}_{2}^{'} < \ldots < \boldsymbol{\theta}_{M}^{'} \end{array}$$

Scattering matrix

$$S_{N}^{M}(\{\theta\},\{\theta'\}) = \\ {}^{out}\langle \theta_{1}^{'},\theta_{2}^{'},\ldots,\theta_{M}^{'}|\theta_{1},\theta_{2},\ldots,\theta_{N}\rangle^{in}$$



Bulk initial state

$$\begin{aligned} &|\theta_1, \theta_2, \dots, \theta_N \rangle^{in} \\ &\theta_1 > \theta_2 > \dots > \theta_N \end{aligned}$$

Form factors

Bulk final state

$$\begin{array}{c} |\boldsymbol{\theta}_{1}^{'},\boldsymbol{\theta}_{2}^{'},\ldots,\boldsymbol{\theta}_{M}^{'}\rangle^{out} \\ \boldsymbol{\theta}_{1}^{'} < \boldsymbol{\theta}_{2}^{'} < \ldots < \boldsymbol{\theta}_{M}^{'} \end{array}$$

Form factors of local operators

$$F_{NM}^{\mathcal{O}}(\{\theta\},\{\theta'\}) = e^{it} \langle \theta_1', \theta_2', \dots, \theta_M' | \mathcal{O}(x, it) | \theta_1, \theta_2, \dots, \theta_N \rangle^{in}$$



Bulk initial state

$$\begin{aligned} & |\theta_1, \theta_2, \dots, \theta_N \rangle^{in} \\ & \theta_1 > \theta_2 > \dots > \theta_N \end{aligned}$$

•

Simplest nontrivial scattering matrix $S_2^2(|\theta_1 - \theta_2|) = S(\theta)$ $\theta > 0$



Unitarity



$$S(\theta)^* = S(-\theta)$$






Scattering matrix: properties (Lorentz invariance)



 $S(\theta) = S(i\pi - \theta)$

Need for analytical continuation in θ

Need for singularity structure in θ

Scattering matrix: properties (Lorentz invariance)











$$< out|T(\Phi(x',it')\mathcal{O}(x,it))|\theta_2,\ldots,\theta_n >$$

$$\begin{array}{c} \hline \text{need: analytic structure of} \\ \hline S \text{ matrix} \end{array} \leftrightarrow \hline \text{reduction formula} \leftrightarrow \hline \begin{array}{c} & \text{known: analytic structure of} \\ \hline \text{correlators} \end{array} \\ \hline < out |\mathcal{O}(x,it)|\theta_1, \theta_2, \dots, \theta_n > = \boxed{2\pi\delta(\theta_1 - \theta)\langle out \setminus \theta | \mathcal{O}(x,it)|\theta_2, \dots, \theta_n >} \\ \hline -Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{-i(i\omega(\theta_1))(it')} \int_{-\infty}^{\infty} dx' e^{-ip(\theta_1)x'} \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2\} \\ \hline < out |T(\Phi(x',it')\mathcal{O}(x,it))|\theta_2, \dots, \theta_n > \end{array}$$

Analytic continuation in θ_1 :

(time reversal) $it \rightarrow -it$ continuation: $\theta \rightarrow i\pi - \theta$

Crossing
$$(x, it) \to (-iT, X) \leftrightarrow (\omega, ip) \to (-iP, \Omega)$$
 continuation $\theta \to \frac{i\pi}{2} - \theta$

Analytic continuation in θ_1 :

(time reversal) $it \rightarrow -it$ continuation: $\theta \rightarrow i\pi - \theta$

Crossing $(x, it) \to (-iT, X) \leftrightarrow (\omega, ip) \to (-iP, \Omega)$ continuation $\theta \to \frac{i\pi}{2} - \theta$

Analytical properties of the correlators: Landau equations

One point function $\langle 0|\Phi(x,it)|0\rangle = G^1(x,it) = G^1$



One point function $\langle 0|\Phi(x,it)|0
angle$



One point function $\langle 0|\Phi(x,it)|0\rangle$



Momentum space formulation $\langle 0|\Phi(x,it)|0\rangle = \int_{-\infty}^{\infty} dx e^{ipx} \int_{-\infty}^{\infty} d(it) e^{i(i\omega)(it)} G^{1}(p,i\omega)$

 $G^1(p,i\omega) = (2\pi)^2 \delta(p) \delta(i\omega) G^1$















$$= \int_{-\infty}^{\infty} dx e^{ipx} \int_{-\infty}^{\infty} d(it) e^{i(i\omega)(it)} G^{1}(p, i\omega)$$



One point function $\langle 0|\Phi(x,it)|0\rangle$



$$= \int_{-\infty}^{\infty} dx e^{ipx} \int_{-\infty}^{\infty} d(it) e^{i(i\omega)(it)} G^{1}(p, i\omega)$$





One point function $\langle 0|\Phi(x,it)|0\rangle$



$$= \int_{-\infty}^{\infty} dx e^{ipx} \int_{-\infty}^{\infty} d(it) e^{i(i\omega)(it)} G^{1}(p, i\omega)$$





 $G^{2}(p_{1}, i\omega_{1}, p_{2}, i\omega_{2}) = (2\pi)^{2}\delta(p_{1} + p_{2})\delta(\omega_{1} + \omega_{2})G^{2}(p, i\omega)$







$$= \int_{-\infty}^{\infty} dx e^{ipx} \int_{-\infty}^{\infty} d(it) e^{i(i\omega)(it)} G^{1}(p, i\omega)$$





 $G^{2}(p_{1}, i\omega_{1}, p_{2}, i\omega_{2}) = (2\pi)^{2}\delta(p_{1} + p_{2})\delta(\omega_{1} + \omega_{2})G^{2}(p, i\omega)$





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$$G^2(p,i\omega) =$$



$$G^{2}(p_{1}, i\omega_{1}, p_{2}, i\omega_{2}) = (2\pi)^{2}\delta(p_{1} + p_{2})\delta(\omega_{1} + \omega_{2})G^{2}(p, i\omega)$$





(x₂,it₂)



(x₁,it)

 $(\mathbf{p}_{2},\mathbf{i}\omega_{2})$

 $(p_1, i\omega_1)$

Correlation functions: higher

Correlation functions: higher



Three point function

Correlation functions: higher



 $|\langle out|\mathcal{O}(x,it)|\theta_1,\theta_2,\ldots,\theta_n\rangle = [2\pi\delta(\theta_1-\theta)\langle out\setminus\theta|\mathcal{O}(x,it)|\theta_2,\ldots,\theta_n\rangle]$

$$< out |\mathcal{O}(x, it)|\theta_1, \theta_2, \dots, \theta_n > = 2\pi \delta(\theta_1 - \theta) \langle out \setminus \theta | \mathcal{O}(x, it)|\theta_2, \dots, \theta_n >$$
$$-Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{-i(i\omega(\theta_1))(it')} \int_{-\infty}^{\infty} dx' e^{-ip(\theta_1)x'} \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2\}$$
$$< out |T(\Phi(x', it')\mathcal{O}(x, it))|\theta_2, \dots, \theta_n >$$

$$< out |\mathcal{O}(x, it)|\theta_1, \theta_2, \dots, \theta_n > = 2\pi \delta(\theta_1 - \theta) \langle out \setminus \theta | \mathcal{O}(x, it)|\theta_2, \dots, \theta_n >$$
$$-Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{-i(i\omega(\theta_1))(it')} \int_{-\infty}^{\infty} dx' e^{-ip(\theta_1)x'} \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2\}$$
$$< out |T(\Phi(x', it')\mathcal{O}(x, it))|\theta_2, \dots, \theta_n >$$

Diagrammatically



$$< out |\mathcal{O}(x, it)|\theta_1, \theta_2, \dots, \theta_n > = 2\pi\delta(\theta_1 - \theta)\langle out \setminus \theta | \mathcal{O}(x, it)|\theta_2, \dots, \theta_n >$$
$$-Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{-i(i\omega(\theta_1))(it')} \int_{-\infty}^{\infty} dx' e^{-ip(\theta_1)x'} \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2\}$$
$$< out |T(\Phi(x', it')\mathcal{O}(x, it))|\theta_2, \dots, \theta_n >$$

Diagrammatically



$$< out |\mathcal{O}(x, it)|\theta_1, \theta_2, \dots, \theta_n > = 2\pi \delta(\theta_1 - \theta) \langle out \setminus \theta | \mathcal{O}(x, it)|\theta_2, \dots, \theta_n >$$
$$-Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{-i(i\omega(\theta_1))(it')} \int_{-\infty}^{\infty} dx' e^{-ip(\theta_1)x'} \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2\}$$
$$< out |T(\Phi(x', it')\mathcal{O}(x, it))|\theta_2, \dots, \theta_n >$$

Diagrammatically



Singularity in correlation function = on mass shell particles



Singularity in correlation function = on mass shell particles



Singularity in correlation function = on mass shell particles



Singularity in correlation function = on mass shell particles



Coleman-Norton interpretation



Singularity in correlation function = on mass shell particles





Coleman-Norton interpretation




Singularity in correlation function = on mass shell particles





Coleman-Norton interpretation







Singularity in correlation function = on mass shell particles



Coleman-Norton interpretation







Cutkosky rules

 $(G^3(u_1))^2$

Singularity in correlation function = on mass shell particles



Coleman-Norton interpretation



Cutkosky rules

 $(G^3(u_1))^2$





 $(G^3(u_1))^2$

Singularity in correlation function = on mass shell particles



Coleman-Norton interpretation







Cutkosky rules

 $(G^3(u_1))^2$

 $(G^3(u_1))^2$

 $(G^{3}(u_{2}))^{4}G^{4}(u_{3})$

Integrability: shifting the trajectories \rightarrow factorization + boostrap

The only nontrivial scattering matrix $S_2^2(|\theta_1 - \theta_2|) = S(\theta)$ $\theta > 0$





Unitarity





Simplest nontrivial solution $S(\theta) = \frac{\sinh \theta - i \sin p}{\sinh \theta + i \sin p}$



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Simplest nontrivial solution $S(\theta) = \frac{\sinh \theta - i \sin p}{\sinh \theta + i \sin p}$ Minimality: all singularity has physical origin: p > 0 end of the story (sinh-Gordon):





for $p = -\frac{2}{3}$ self-fusion: Lee-Yang, for generic p sine-Gordon $B_2, B_3, \ldots, B_n, s, \overline{s}$

Permutation $F_{n}^{\mathcal{O}}(\theta_{1}, \dots, \theta_{i}, \theta_{i+1}, \dots, \theta_{n}) =$ $S(\theta_{i} - \theta_{i+1})F_{n}^{\mathcal{O}}(\theta_{1}, \dots, \theta_{i+1}, \theta_{i}, \dots, \theta_{n})$









Periodicity

$$F_n^{\mathcal{O}}(\theta_1, \dots, \theta_{n-1}, \theta_n) =$$

 $F_n^{\mathcal{O}}(\theta_n + 2\pi i, \theta_1, \dots, \theta_{n-1})$





Kinematical singularities

$$i^{\text{res}}_{\theta=\theta'}F_{n+2}^{\mathcal{O}}(\theta+i\pi,\theta',\theta_1,\ldots,\theta_n) = \left(1-\prod_{i=1}^n S(\theta-\theta_i)\right)F_n^{\mathcal{O}}(\theta_1,\ldots,\theta_n)$$





Dynamical singuralities

$$i \operatorname{res}_{\theta = \theta'} F_{n+2}^{\mathcal{O}}(\theta + iu, \theta' - iu, \theta_1, \dots, \theta_n) = \Gamma F_{n+1}^{\mathcal{O}}(\theta, \theta_1, \dots, \theta_n)$$





Form factor solutions for generic p: sinh-Gordon, Lee-Yang, B_1 form factor in sine-Gordon

Boundary Hilbert space

Boundary initial state

$$| heta_1, heta_2,\ldots, heta_N; E
angle_B^{in}$$

 $heta_1> heta_2>\ldots> heta_N>\mathbf{0}$



Boundary Hilbert space

Boundary final state

 $\begin{aligned} & |\boldsymbol{\theta}_{1}^{'},\boldsymbol{\theta}_{2}^{'},\ldots,\boldsymbol{\theta}_{M}^{'};\boldsymbol{E}^{'}\rangle_{B}^{out} \\ & \boldsymbol{\theta}_{1}^{'} < \boldsymbol{\theta}_{2}^{'} < \ldots < \boldsymbol{\theta}_{M}^{'} < \boldsymbol{0} \end{aligned}$

Boundary initial state

 $| heta_1, heta_2, \dots, heta_N; E
angle_B^{in}$ $heta_1 > heta_2 > \dots > heta_N > \mathbf{0}$



Reflection matrix

Boundary final state

$$\theta_{1}^{'} < \theta_{2}^{'} < \dots < \theta_{M}^{'}; \frac{E^{'}}{B} \\ \theta_{1}^{'} < \theta_{2}^{'} < \dots < \theta_{M}^{'} < \mathbf{0}$$

Reflection matrix

$$\underset{B}{\overset{out}{B}} \begin{pmatrix} \theta'_{1}, \theta'_{2}, \dots, \theta'_{M}; E' | \theta_{1}, \theta_{2}, \dots, \theta_{N}; E \rangle_{B}^{in} \end{pmatrix} =$$

Boundary initial state

$$|\theta_1, \theta_2, \dots, \theta_N; E\rangle_B^{in}$$

 $\theta_1 > \theta_2 > \dots > \theta_N > 0$



Analytic structure?

Boundary form factors

Boundary initial state

$$|\theta_1, \theta_2, \dots, \theta_N; E\rangle_B^{in}$$

 $\theta_1 > \theta_2 > \dots > \theta_N > 0$



Analytic structure?





$$\begin{array}{c|c} \text{need: analytic structure of} \\ \hline R \text{ matrix} \end{array} \leftrightarrow \hline \text{boundary reduction formula} \leftrightarrow \hline \text{analytic structure of} \\ \hline boundary \text{ correlators} \end{array}$$

$$B < out |\mathcal{O}(it)|\theta_1, \theta_2, \dots, \theta_n >_B = 2\pi\delta(\theta_1 - \theta) B < out \setminus \theta |\mathcal{O}(it)|\theta_2, \dots, \theta_n >_B \\ \hline -2Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{i(i\omega(\theta_1))(it')} \int_{-\infty}^{0} dx' \cos(p(\theta_1)x') \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2 + \delta(x')\partial_{x'}\}$$

$$\begin{array}{c} \begin{array}{c} \text{need: analytic structure of} \\ \hline R \text{ matrix} \end{array} \leftrightarrow \begin{array}{c} \begin{array}{c} \text{boundary reduction formula} \end{array} \leftrightarrow \begin{array}{c} \begin{array}{c} \text{analytic structure of} \\ \hline \text{boundary correlators} \end{array} \end{array}$$

$$\begin{array}{c} B < out |\mathcal{O}(it)|\theta_1, \theta_2, \dots, \theta_n >_B \end{array} = \left[2\pi\delta(\theta_1 - \theta)_B < out \setminus \theta |\mathcal{O}(it)|\theta_2, \dots, \theta_n >_B \right] \end{array}$$

$$\begin{array}{c} -2Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{i(i\omega(\theta_1))(it')} \int_{-\infty}^{0} dx' \cos(p(\theta_1)x') \{-\partial_{it'}^2 - \partial_{x'}^2 + m^2 + \delta(x')\partial_{x'}\} \end{array}$$

 $B < out|T(\Phi(x', it')\mathcal{O}(it))|\theta_2, \dots, \theta_n >_B$

$$\begin{array}{c} \begin{array}{c} \text{need: analytic structure of} \\ \hline R \text{ matrix} \end{array} \leftrightarrow \begin{array}{c} \begin{array}{c} \text{boundary reduction formula} \end{array} \leftrightarrow \begin{array}{c} \begin{array}{c} \text{analytic structure of} \\ \hline \text{boundary correlators} \end{array} \end{array}$$

$$\begin{array}{c} B < out |\mathcal{O}(it)|\theta_1, \theta_2, \dots, \theta_n >_B \end{array} = \left[2\pi\delta(\theta_1 - \theta)_B < out \setminus \theta |\mathcal{O}(it)|\theta_2, \dots, \theta_n >_B \right] \end{array}$$

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$$B < out|T(\Phi(x', it')\mathcal{O}(it))|\theta_2, \dots, \theta_n >_B$$

Analytic continuation in θ_1 :

(left/right matrix element)
$$\rightarrow$$
 continuation: $\theta \rightarrow i\pi + \theta$

Crossing ?

$$\begin{array}{c} \begin{array}{c} \text{need: analytic structure of} \\ \hline R \text{ matrix} \end{array} \leftrightarrow \begin{array}{c} \text{boundary reduction formula} \end{array} \leftrightarrow \begin{array}{c} \text{analytic structure of} \\ \hline boundary \text{ correlators} \end{array} \end{array}$$

$$\begin{array}{c} B < out |\mathcal{O}(it)|_{\theta_1}, \theta_2, \dots, \theta_n >_B \end{array} = \left[2\pi\delta(\theta_1 - \theta)_B < out \setminus \theta |\mathcal{O}(it)|_{\theta_2}, \dots, \theta_n >_B \right] \\ \hline -2Z^{-1/2} \int_{-\infty}^{\infty} d(it') e^{i(i\omega(\theta_1))(it')} \int_{-\infty}^{0} dx' \cos(p(\theta_1)x') \{ -\partial_{it'}^2 - \partial_{x'}^2 + m^2 + \delta(x')\partial_{x'} \} \end{array}$$

$$B < out|T(\Phi(x', it')\mathcal{O}(it))|\theta_2, \ldots, \theta_n >_B$$

Analytic continuation in θ_1 :

(left/right matrix element) \rightarrow continuation: $\theta \rightarrow i\pi + \theta$

Crossing ?

Analytical properties of the correlators: Landau equations
One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B} = G_{B}^{1}(x,it) = G_{B}^{1}(x)$



One point function ${}_{B}\langle 0|\Phi(x,it)|0\rangle_{B} = G^{1}_{B}(x,it) = G^{1}_{B}(x)$



One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$



One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$



Momentum space formulation $_{B}\langle 0|\Phi(x,it)|0\rangle_{B} = \int_{-\infty}^{\infty} dx e^{ipx} \int_{-\infty}^{\infty} dt e^{i(i\omega)(it)}G^{1}_{B}(p,i\omega)$

 $G^{1}_{B}(p,i\omega) = (2\pi)\delta(i\omega)G^{2}(p,i\omega)G^{1}_{B}(p)$





One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$







(p, iω)







One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$







Two point function

 ${}_{B}\langle 0|T(\Phi(x_{1},it_{1})\Phi(x_{2},it_{2}))|0\rangle_{B} = G_{B}^{2}(x_{1},x_{2},it_{1}-it_{2})$



One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$







Two point function

 ${}_{B}\langle 0|T(\Phi(x_{1},it_{1})\Phi(x_{2},it_{2}))|0\rangle_{B} = G_{B}^{2}(x_{1},x_{2},it_{1}-it_{2})$





One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$







Two point function

 $_{B}\langle 0|T(\Phi(x_{1},it_{1})\Phi(x_{2},it_{2}))|0\rangle_{B} = G_{B}^{2}(x_{1},x_{2},it_{1}-it_{2})$



 $G_{B}^{2}(p_{1}, i\omega_{1}, p_{2}, i\omega_{2}) = (2\pi)\delta(i\omega_{1} + i\omega_{2})G_{B}^{2}(p_{1}, p_{2}, i\omega_{1})$

One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$







Two point function

$${}_{B}\langle 0|T(\Phi(x_{1},it_{1})\Phi(x_{2},it_{2}))|0\rangle_{B} = G^{2}_{B}(x_{1},x_{2},it_{1}-it_{2})$$

 $G_{B}^{2}(p_{1}, i\omega_{1}, p_{2}, i\omega_{2}) = (2\pi)\delta(i\omega_{1} + i\omega_{2})G_{B}^{2}(p_{1}, p_{2}, i\omega_{1})$







 $(p_1, i\omega_1)$





Two point function

 $_{B}\langle 0|T(\Phi(x_{1},it_{1})\Phi(x_{2},it_{2}))|0\rangle_{B} = G_{B}^{2}(x_{1},x_{2},it_{1}-it_{2})$

 $G_{B}^{2}(p_{1}, i\omega_{1}, p_{2}, i\omega_{2}) = (2\pi)\delta(i\omega_{1} + i\omega_{2})G_{B}^{2}(p_{1}, p_{2}, i\omega_{1}) (p_{2}i\omega_{2})$





 $\boxed{B < \theta_1, \theta_2, \dots, \theta_n | \mathcal{O}(t) | in >_B} = \boxed{\delta(\theta_1 - \theta)_B < \theta_2, \dots, \theta_n | \mathcal{O}(t) | in \setminus \theta >_B}$

$$B < \theta_1, \theta_2, \dots, \theta_n |\mathcal{O}(t)| in >_B = \delta(\theta_1 - \theta)_B < \theta_2, \dots, \theta_n |\mathcal{O}(t)| in \setminus \theta >_B$$
$$i2Z^{-1/2} \int_{-\infty}^{\infty} dt' e^{-i\omega(\theta_1)t'} \int_{-\infty}^{0} dx' \cos(p(\theta_1)x') \{\partial_{t'}^2 - \partial_{x'}^2 + m^2 + \delta(x')\partial_{x'}\}$$

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$$B < \theta_2, \ldots, \theta_n | T(\Phi(x', t') \mathcal{O}(t)) | in > B$$

$$B < \theta_1, \theta_2, \dots, \theta_n | \mathcal{O}(t) | in >_B = \delta(\theta_1 - \theta) | B < \theta_2, \dots, \theta_n | \mathcal{O}(t) | in \setminus \theta >_B$$
$$i2Z^{-1/2} \int_{-\infty}^{\infty} dt' e^{-i\omega(\theta_1)t'} \int_{-\infty}^{0} dx' \cos(p(\theta_1)x') \{\partial_{t'}^2 - \partial_{x'}^2 + m^2 + \delta(x')\partial_{x'}\}$$

$$B < \theta_2, \ldots, \theta_n | T(\Phi(x', t') \mathcal{O}(t)) | in >_B$$

Diagrammatically



$$B < \theta_1, \theta_2, \dots, \theta_n | \mathcal{O}(t) | in >_B = \delta(\theta_1 - \theta)_B < \theta_2, \dots, \theta_n | \mathcal{O}(t) | in \setminus \theta >_B$$
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Diagrammatically



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$$B < \theta_2, \ldots, \theta_n | T(\Phi(x', t') \mathcal{O}(t)) | in >_B$$

Diagrammatically



Singularity in correlation function = on mass shell particles



Singularity in correlation function = on mass shell particles



Singularity in correlation function = on mass shell particles







Singularity in correlation function = on mass shell particles





Coleman-Norton type interpretation



Singularity in correlation function = on mass shell particles





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Singularity in correlation function = on mass shell particles





Coleman-Norton type interpretation







Cutkosky type rules

 $(G^1_{\underline{B}}(u_1))^2$

Singularity in correlation function = on mass shell particles





Coleman-Norton type interpretation





Cutkosky type rules

 $(G^1_{B}(u_1))^2$

 $(G^3(u_1))G^1_{B}(0)$

Singularity in correlation function = on mass shell particles





Coleman-Norton type interpretation



Cutkosky type rules

 $(G^1_{\underline{B}}(u_1))^2$





 $(G^3(u_1))G^1_B(0)$

 $(G^3(u_2))^2 G^2_{\mathbf{R}}(u_3)$

Integrability: shifting the trajectories \rightarrow factorization + boostrap

The only nontrivial is the 1pt reflection matrix $R(|\theta_1|) = R(\theta)$ $\theta > 0$





Unitarity









Minimality: all singularity has physical origin:


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Minimality: all singularity has physical origin:

Bootstrap









 $\begin{array}{l}
 Reflection \\
 F_n^{\mathcal{O}}(\theta_1, \dots, \theta_{n-1}, \theta_n) = \\
 R(\theta_n) F_n^{\mathcal{O}}(\theta_1, \dots, \theta_{n-1}, -\theta_n)
\end{array}$



 $\begin{array}{l} \text{Reflection} \\ F_n^{\mathcal{O}}(\theta_1, \dots, \theta_{n-1}, \theta_n) = \\ R(\theta_n) F_n^{\mathcal{O}}(\theta_1, \dots, \theta_{n-1}, -\theta_n) \end{array}$







 θ_n

θ_n



Boundary periodicity

$$F_n^{\mathcal{O}}(\theta_1, \theta_2, \dots, \theta_n) =$$

 $R(i\pi - \theta_1)F_n^{\mathcal{O}}(2i\pi - \theta_1, \theta_2, \dots, \theta_n)$



 θ_n

$$\begin{array}{l} \text{Kinematical singularities} \\ i \operatorname{res}_{\theta = \theta'} F_{n+2}^{\mathcal{O}}(\theta + i\pi, \theta', \theta_1, \dots, \theta_n) = \\ \left(1 - \prod_{i=1}^n S(\theta - \theta_i) S(\theta + \theta_i)\right) F_n^{\mathcal{O}}(\theta_1, \dots, \theta_n) \\ i \operatorname{res}_{\theta = \theta'} F_{n+2}^{\mathcal{O}}(-\theta + i\pi, \theta', \theta_1, \dots, \theta_n) = \\ \left(R(\theta) - \prod_{i=1}^n S(\theta - \theta_i) R(\theta) S(\theta + \theta_i)\right) F_n^{\mathcal{O}}(\theta_1, \dots, \theta_n)
\end{array}$$





Bulk dynamical singuralities

$$i \operatorname{res}_{\theta = \theta'} F_{n+2}^{\mathcal{O}}(\theta + iu, \theta' - iu, \theta_1, \dots, \theta_n) = \Gamma F_{n+1}^{\mathcal{O}}(\theta, \theta_1, \dots, \theta_n)$$





Boundary dynamical singuralities

$$i \operatorname{res}_{\theta = iu} F_{n+1}^{\mathcal{O}}(\theta_1, \dots, \theta_n, \theta) = g \tilde{F}^{\mathcal{O}}(\theta_1, \dots, \theta_n)$$



One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B} = G^{1}_{B}(x,it) = G^{1}_{B}(x)$



Path integral representation:

 ${}_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$

 $= Z^{-1}{}_{B}\langle 0|U_{B}(\infty,it)\Phi(x,it)U_{B}(it,-\infty)|0\rangle_{B}$



Path integral representation:

 $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$

 $= Z^{-1}{}_{B}\langle 0|U_{B}(\infty,it)\Phi(x,it)U_{B}(it,-\infty)|0\rangle_{B}$

 $=\int \mathcal{D}\Phi \Phi(x,it)e^{-S[\Phi(x,it)]}$



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 $_{B}\langle 0|\Phi(x,it)|0\rangle_{B}$

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 $=\int \mathcal{D}\Phi \Phi(x,it)e^{-S[\Phi(x,it)]}$

 $= \int \mathcal{D}\Phi \,\Phi(-iT,X) e^{-S[\Phi(-iT,X)]}$









One point function $_{B}\langle 0|\Phi(x,it)|0\rangle_{B} = G_{B}^{1}(x,it) = G_{B}^{1}(x)$



One point function $\langle B | \Phi(-iT, X) | 0 \rangle^{in} = G^1_B(-iT, X) = G^1_B(-iT)$



One point function $\langle B | \Phi(-iT, X) | 0 \rangle^{in} = G^1_B(-iT, X) = G^1_B(-iT)$



One point function $\langle B | \Phi(-iT, X) | 0 \rangle^{in} = G^1_B(-iT, X) = G^1_B(-iT)$

Momentum space: $\langle B | \Phi(-iT, X) | 0 \rangle = \int_{-\infty}^{\infty} dx e^{iPX} \int_{-\infty}^{\infty} d(iT) e^{i(-i\Omega)(-iT)} G_B^1(P, -i\Omega)$

(x=-iT,it=x

 $G^{1}(P,-i\Omega) = (2\pi)\delta(P)G^{2}(P,-i\Omega)G^{1}_{B}(-i\Omega)$



One point function $\langle \mathbf{B} | \Phi(-iT, X) | 0 \rangle^{in} = G^1_B(-iT, X) = G^1_B(-iT)$



$$= \int_{-\infty}^{\infty} dX e^{iPX} \int_{-\infty}^{\infty} d(-iT) e^{i(-i\Omega)(-iT)} G^{1}_{B}(P, -i\Omega)$$



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(x=-iT,it=x)

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(x=-iT,it=x)



One point function $\langle B | \Phi(-iT, X) | 0 \rangle^{in} = G^1_B(-iT, X) = G^1_B(-iT)$

$$= \int_{-\infty}^{\infty} dX e^{iPX} \int_{-\infty}^{\infty} d(-iT) e^{i(-i\Omega)(-iT)} G^{1}_{B}(P,-i\Omega)$$

Two point function

 $\langle \mathbf{B} | T(\Phi(-iT_1, X_1) \Phi(-iT_2, X_2)) | 0 \rangle = G_{\mathbf{B}}^2(-iT_1, -iT_2, X_1 - X_2)$ (iT₁, X₁)

 $G_{B}^{2}(P_{1}, i\Omega_{1}, P_{2}, i\Omega_{2}) = (2\pi)\delta(P_{1} + P_{2})G_{B}^{2}(P_{1}, i\Omega_{1}, i\Omega_{2})$







One point function
$$\langle B | \Phi(-iT, X) | 0 \rangle^{in} = G^1_B(-iT, X) = G^1_B(-iT)$$

$$= \int_{-\infty}^{\infty} dX e^{iPX} \int_{-\infty}^{\infty} d(-iT) e^{i(-i\Omega)(-iT)} G^{1}_{B}(P,-i\Omega)$$

Two point function

 $\langle \mathbf{B} | T(\Phi(-iT_1, X_1) \Phi(-iT_2, X_2)) | 0 \rangle = G_{\mathbf{B}}^2(-iT_1, -iT_2, X_1 - X_2)$

 $G_{B}^{2}(P_{1}, i\Omega_{1}, P_{2}, i\Omega_{2}) = (2\pi)\delta(P_{1} + P_{2})G_{B}^{2}(P_{1}, i\Omega_{1}, i\Omega_{2})$



(iT, X)





Boundary state: definition

Boundary state: definition

One point function

$$\langle \mathbf{B} | \Phi(iT, X) | \mathbf{0} \rangle = G_{\mathbf{B}}^{1}(iT, X) = {}_{\mathbf{B}} \langle \mathbf{0} | \Phi(x, it) | \mathbf{0} \rangle_{\mathbf{B}} = G_{\mathbf{B}}^{1}(x, it)$$


Boundary state: definition

One point function

$$\langle \mathbf{B} | \Phi(iT, X) | \mathbf{0} \rangle = G_{\mathbf{B}}^{1}(iT, X) = {}_{\mathbf{B}} \langle \mathbf{0} | \Phi(x, it) | \mathbf{0} \rangle_{\mathbf{B}} = G_{\mathbf{B}}^{1}(x, it)$$



Two point function

 $\langle \mathbf{B} | T(\Phi(iT_1, X_1) \Phi(iT_2, X_2)) | \mathbf{0} \rangle = G_{\mathbf{B}}^2(iT_1, iT_2, X_1 - X_2) = \mathbf{B} \langle \mathbf{0} | T(\Phi(x_1, it_1) \Phi(x_2, it_2)) | \mathbf{0} \rangle_{\mathbf{B}}$





One particle contribution

 $\langle \mathbf{B} | \Theta \rangle^{in} = K^1 2\pi \delta(P(\Theta))$



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Two particle contribution

$$\langle \mathbf{B} | \Theta_1, \Theta_2 \rangle^{in} = K^2(\Theta_1) 2\pi \delta(P(\Theta_1) + P(\Theta_2))$$



One particle contribution

 $\langle \mathbf{B} | \Theta \rangle^{in} = K^1 2\pi \delta(P(\Theta))$



n particle contribution

$$\langle \mathbf{B} | \Theta_1, \Theta_2, \dots, \Theta_n \rangle^{in} =$$

 $K^n(\Theta_1,\ldots,\Theta_n) 2\pi \delta(\sum_i P(\Theta_i))$



Two particle contribution

 $\langle \mathbf{B} | \Theta_1, \Theta_2 \rangle^{in} = K^2(\Theta_1) 2\pi \delta(P(\Theta_1) + P(\Theta_2))$



One particle contribution

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$$\langle \mathbf{B} | \Theta_1, \Theta_2, \dots, \Theta_n \rangle^{in} =$$

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Two particle contribution

$$\langle B|\Theta_1,\Theta_2\rangle^{in} = K^2(\Theta_1)2\pi\delta(P(\Theta_1) + Connection to correlators?) + Connection to correlators? Crossed channel reduction formula$$



 $< \theta_1, \theta_2, \dots, \theta_n | B > =$

$$\langle \theta_1, \theta_2, \dots, \theta_n | B \rangle =$$

$$\int_{-\infty}^{\infty} dX' \int_{-\infty}^{0} dT' e^{-i\Omega(\theta_1)T' + iP(\theta_1)X'} \{-\partial_{iT'}^2 - \partial_{X'}^2 + m^2 - \delta(iT')(\partial_{iT'} + i\Omega)\}$$

$$\langle \theta_1, \theta_2, \dots, \theta_n | B \rangle =$$

$$| iZ^{-1/2} \int_{-\infty}^{\infty} dX' \int_{-\infty}^{0} dT' e^{-i\Omega(\theta_1)T' + iP(\theta_1)X'} \{ -\partial_{iT'}^2 - \partial_{X'}^2 + m^2 - \delta(iT')(\partial_{iT'} + i\Omega) \}$$

$$|< heta_2,\ldots, heta_n|\Phi(x^{'},t^{'})|B>|$$

$$\langle \theta_1, \theta_2, \dots, \theta_n | B \rangle =$$

$$iZ^{-1/2} \int_{-\infty}^{\infty} dX' \int_{-\infty}^{0} dT' e^{-i\Omega(\theta_1)T' + iP(\theta_1)X'} \{-\partial_{iT'}^2 - \partial_{X'}^2 + m^2 - \delta(iT')(\partial_{iT'} + i\Omega)\}$$

$$|< heta_2,\ldots, heta_n|\Phi(x',t')|B>|$$

Diagrammatically



$$\langle \theta_1, \theta_2, \dots, \theta_n | B \rangle =$$

$$iZ^{-1/2} \int_{-\infty}^{\infty} dX' \int_{-\infty}^{0} dT' e^{-i\Omega(\theta_1)T' + iP(\theta_1)X'} \{-\partial_{iT'}^2 - \partial_{X'}^2 + m^2 - \delta(iT')(\partial_{iT'} + i\Omega)\}$$

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Diagrammatically



$$\langle \theta_1, \theta_2, \dots, \theta_n | B \rangle =$$

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$$|< heta_2,\ldots, heta_n|\Phi(x',t')|B>|$$

Diagrammatically



Crossing: Boundary state, reflection

Crossing: Boundary state, reflection

One particle contribution \leftrightarrow One particle emission (reflection)



Crossing: Boundary state, reflection

One particle contribution ↔ One particle emission (reflection)





Two particle contribution \leftrightarrow One particle reflection





Reflection factor is related to the twopoint function

 ${}_{\boldsymbol{B}}\langle 0|T(\Phi(x_1,it_1)\Phi(x_2,it_2))|0\rangle_{\boldsymbol{B}}$



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 $_{\boldsymbol{B}}\langle 0|T(\Phi(x_1,it_1)\Phi(x_2,it_2))|0\rangle_{\boldsymbol{B}}$

Choosing large separation in $t_1 - t_2$

 $\lim_{(t_1-t_2)\to\infty} {}_{B}\langle 0|T(\Phi(x_1,it_1)\Phi(x_2,it_2))|0\rangle_{B} =$



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Choosing large separation in $t_1 - t_2$

$$\lim_{(t_1-t_2)\to\infty} {}_{\mathbf{B}}\langle 0|T(\Phi(x_1,it_1)\Phi(x_2,it_2))|0\rangle_{\mathbf{B}} =$$

 ${}_{B}\langle 0|\Phi(x_{2},it_{2})|0\rangle_{BB}\langle 0|\Phi(x_{2},it_{2})|0\rangle_{B}$



Reflection factor is related to the twopoint function

 $_{\boldsymbol{B}}\langle 0|T(\Phi(x_1,it_1)\Phi(x_2,it_2))|0\rangle_{\boldsymbol{B}}$

Choosing large separation in $t_1 - t_2$

 $\lim_{(t_1-t_2)\to\infty} {}_{\boldsymbol{B}}\langle 0|T(\Phi(x_1,it_1)\Phi(x_2,it_2))|0\rangle_{\boldsymbol{B}} =$

 $_{B}\langle 0|\Phi(x_{2},it_{2})|0\rangle_{BB}\langle 0|\Phi(x_{2},it_{2})|0\rangle_{B}$

Reflection factor
$$\left[\left. R_1^1(\theta) \right|_{\theta \to \frac{i\pi}{2}} = \frac{g^2}{2} \frac{1}{i(\theta - \frac{i\pi}{2})} = \frac{g^2}{2} \frac{1}{i\Theta} = K^2(\Theta) \right]$$

'Emission' factor $R^1 = \frac{g}{2} = K^1$

Example
$$R(\theta) = \frac{\binom{1}{2} \binom{1-\frac{p}{2}}{\pi} \frac{\binom{i\eta p}{\pi} - \frac{1}{2}}{\binom{\frac{1}{2} - \frac{p}{2}}{\pi} \frac{\binom{ip \theta}{\pi} - \frac{1}{2}}{\binom{\frac{ip \theta}{\pi} - \frac{1}{2}}{(\frac{ip \theta}{\pi} + \frac{1}{2})}}$$
; $R^1 = \sqrt{\cot \frac{\pi B}{4} \cot \frac{\pi (1-B)}{4}} \tan \frac{\eta B}{2} \tanh \frac{\theta B}{2}$



One point function:

 $\langle \mathbf{B} | \Phi(-iT, X) | \mathbf{0} \rangle =$



One point function: Vacuum+

 $\langle B | \Phi(-iT, X) | 0 \rangle = \langle B | 0 \rangle \times \\ \langle 0 | \Phi(-iT, X) | 0 \rangle +$ (x = -iT, it = x) (x = -iT, it = x)





Ground state energy: partition function

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Ground state energy from the partition function

$$\lim_{R\to\infty} Z(L,R) = \lim_{R\to\infty} Tr(e^{-H_{\mathbf{B}}(L)R})$$

$$= e^{-E_0(L)R} \to E_0(L) = -\lim_{R \to \infty} R^{-1} \log Z(L,R)$$



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Alternative calculation from the crossed channel

$$\lim_{R\to\infty} Z(L,R) =$$

 $\lim_{R\to\infty} \langle \mathbf{B}_0 | e^{-H(R)L} | \mathbf{B}_L \rangle$





Partition function:

$$Z(L,R) = \langle \mathbf{B}_0 | e^{-H(R)L} | \mathbf{B}_L \rangle$$









Ground state: Casimir energy

Ground state: Casimir energy

Nonvanishing *g*:

 $E_0(L) = -2mK_0^1 K_1^{1*} e^{-mL} +$


Ground state: Casimir energy



Ground state: Casimir energy



free theory

$$E_{0L}^{0}(L) = \int_{-\infty}^{\infty} \frac{d\theta}{2\pi} m \cosh\theta \log(1 + K_0^2(\theta) K_L^{2*}(\theta) e^{-2m \cosh\theta L})$$