

# Does Planck length challenge non-relativistic quantum mechanics of large masses?

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- 1 Does Planck scale require relativistic motion?
- 2 Nonrelativistic plane wave versus Planckian scale
- 3 Large mass non-relativistic wave function
- 4 Ignorable effects per atoms accumulate
- 5 Is centre-of-mass  $x$  of  $10^{23}$  atoms observable at all?
- 6 Facts and questions
- 7 Closing remarks

# Does Planck scale require relativistic motion?

- Space-time continuum is likely to break down at  $\ell_{\text{Pl}}$ .
- Hence Planck scale puts a limit to standard physics.
- Early Big Bang high energies hit the Planck scale.
- At lower energies than that, we remain on the safe side.
- Which is not quite true.

When quantum mechanics enters the Planck scale:

$$\underbrace{\lambda_{\text{deBrogie}}}_{\frac{2\pi\hbar}{mv} \sqrt{1-v^2/c^2}} \sim \underbrace{\ell_{\text{Planck}}}_{\sqrt{\frac{\hbar G}{c^3}}}$$

$\lambda_{\text{dB}}$  sinks to  $\ell_{\text{Pl}}$  in two ways:

- Relativistic way: when elementary particles velocity  $v$  closes  $c$ .
- Non-relativistic: *when mass grows macroscopic.*

# Nonrelativistic plane wave versus Planckian scale

Non-relativistic plane wave,  $v \ll c$ :

$$\begin{aligned}\Psi(x, t) &= \exp(-iEt/\hbar + ipx/\hbar) \\ &= \exp\left(-i\frac{mv^2}{2\hbar}t + i\frac{mv}{\hbar}x\right) \\ &= \exp\left(-2\pi i\frac{t}{\tau} + 2\pi i\frac{x}{\lambda}\right)\end{aligned}$$

$$\text{Periodicity in } t \text{ and } x: \begin{cases} \tau = (4\pi\hbar/mv^2) \\ \lambda = (2\pi\hbar/mv) \end{cases}$$

$\Psi(x, t)$  is legitimate as long as  $\tau \gg \tau_{\text{Pl}}$  and  $\lambda \gg \ell_{\text{Pl}}$ .

$$\text{Planck time and length: } \begin{cases} \tau_{\text{Pl}} = \sqrt{\hbar G/c^5} \sim 10^{-43} \text{ s} \\ \ell_{\text{Pl}} = \sqrt{\hbar G/c^3} \sim 10^{-33} \text{ cm} \end{cases}$$

For atomic  $m$  that's the case:  $\tau/\tau_{\text{Pl}} \gg 10^{18}$  and  $\lambda/\ell_{\text{Pl}} \gg 10^{18}$ .  
But larger  $m$  will push  $\Psi(x, t)$  towards the Planckian scales.

# Large mass non-relativistic wave function

- $\tau = \frac{2\pi\hbar}{mv^2/2}$  sinks to  $\tau_{\text{Pl}} \sim 10^{-43}\text{s}$  if  $mv^2 \sim m_{\text{Pl}}c^2$ , i.e.:

$$m \sim \frac{c^2}{v^2} m_{\text{Pl}} \sim \frac{c^2}{v^2} 10^{-5}\text{g}.$$

- $\lambda = \frac{2\pi\hbar}{mv}$  sinks to  $\ell_{\text{Pl}} \sim 10^{-33}\text{cm}$  if  $mv \sim m_{\text{Pl}}c$ , i.e.:

$$m \sim \frac{c}{v} m_{\text{Pl}} \sim \frac{c}{v} 10^{-5}\text{g}.$$

With growing  $m$ , non-relativistic  $\Psi(x, t)$  becomes illegitimate.  
The bell rings for spatial periodicity first.

Example I, free motion:

$$m = 10\text{g}, v = 10\text{km/s}, \Rightarrow \lambda = (2\pi\hbar/mv) \sim \ell_{\text{Pl}}$$

Example II, rigid body elastic vibration mode:

$$m = 10\text{kg}, \omega = 100\text{kHz}, a(\text{mplitude}) = 10^{-2}\text{cm} \Rightarrow \lambda = (2\pi\hbar/m a \omega) \sim \ell_{\text{Pl}}$$

# Ignorable effects per atoms accumulate

Suppose a global constant “uncertainty”:

$$x \Rightarrow x + u.$$

$|u| \sim \ell_{\text{Pl}}$  is space's error/bluriness/foaminess/fluctuation.

Many-particle state:

$$\Psi(x_1, x_2, \dots, x_{10^{23}}) \Rightarrow \Psi(x_1 + u, x_2 + u, \dots, x_{10^{23}} + u).$$

$u$  is irrelevant for non-relativistic individual particles, but its *effect accumulates* for the many-particle c.o.m.  $x$ .

Best seen in momentum representation:

$$\tilde{\Psi}(p_1, p_2, \dots, p_{10^{23}}) \Rightarrow \exp\left(\frac{i}{\hbar} u \underbrace{(p_1 + p_2 + \dots + p_{10^{23}})}_P\right) \tilde{\Psi}(p_1, p_2, \dots, p_{10^{23}})$$

C.o.m. reduced state decoheres in momentum if  $u$  is stochastic:

$$\rho(P, P') \Rightarrow \exp\left(\frac{i}{\hbar} u (P - P')\right) \rho(P, P') \Rightarrow \exp\left(-\frac{\ell_{\text{Pl}}^2}{2\hbar^2} (P - P')^2\right) \rho(P, P')$$

# Is c.o.m. operator $x$ of $10^{23}$ atoms observable?

Sure, it is!

In quantum optomechanics, magnetomechanics:

Spatial motion of suspended, flexibly located, levitated or trapped macroobjects is controlled in their quantum regime.

E.g.: Each and every photon in mirror-optomechanics interacts with the mirror as a whole.

Masses are still much less than those requested for  $\lambda \sim \ell_{\text{Pl}}$ .  
But much larger than before 20 years we believed in.

# Facts and questions

- We don't need extreme high energies to explore Planck scale.
- Nonrelativistic QM *of massive d.o.f.* does explore it.
- And breaks down there. What way, we don't yet know.
- Plausible: space-time “uncertainty” yields noise/decoherence.
  - Holographic noise? (Hogan [Genovese's talk])
  - G-related decoherence of massive d.o.f. ? (D-Penrose)
- Breakdown depends on spectrum of “uncertainty”  $u$ .
- Can come much earlier than for global static  $u$ .

D-Penrose: nonrelativistic

- Can it be the non-relativistic footprint of Planck scale “uncertainties”?
- If it decoheres massive d.o.f. before their  $\lambda_{\text{dB}}$  sinks to  $\ell_{\text{Pl}}$ ?



# Closing remarks

The knowledge that

i) *Planckian “uncertainty” of space-time accumulates for large non-relativistic objects,*

ii) *we might therefore study Planckian footprints in the lab non-relativistically,*

has been *implicit* in various works (Károlyházy, D., Penrose, Hogan, Bekenstein ...), all using *sophisticated arguments*.

What I'm adding is *explicit and elementary evidence*.

Remember:

$$\lambda_{\text{dB}} = \frac{2\pi\hbar}{10g \times 10\text{km/s}} = 4.2 \times 10^{-33} \text{cm} \sim \ell_{\text{Pl}}.$$