Earliest stochastic Schrödinger equations from foundations

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Early Motivations 1970-1980's

Interpretation of ψ is statistical. Sudden 'one-shot' collapse $\psi \to \psi_n$ is central.

- If collapse takes time?
- Hunt for a math model (Pearle, Gisin, Diosi)
- New physics?

1-Shot Non-Selective Measurement, Decoherence

Measurement of \hat{A} , pre-measurement state $\hat{\rho}$, post-measurement state, decoherence: $\hat{A} = \sum_n A_n \hat{P}_n$; $\sum_n \hat{P}_n = \hat{I}$, $\hat{P}_n \hat{P}_m = \delta_{nm} \hat{I}$

$$\hat{\rho} \rightarrow \sum_{n} \hat{P}_{n} \hat{\rho} \hat{P}_{n}$$

Off-diagonal elements become zero: Decoherence.

Example:
$$\hat{A} = \hat{\sigma}_Z = |\uparrow\rangle\langle\uparrow| - |\downarrow\rangle\langle\downarrow|, \ \hat{P}_{\uparrow} = |\uparrow\rangle\langle\uparrow|, \ \hat{P}_{\downarrow} = |\downarrow\rangle\langle\downarrow|,$$

$$\hat{\rho} = \rho_{\uparrow\uparrow} \left|\uparrow\right\rangle \left\langle\uparrow\right| + \rho_{\downarrow\downarrow} \left|\downarrow\right\rangle \left\langle\downarrow\right| + \rho_{\uparrow\downarrow} \left|\uparrow\right\rangle \left\langle\downarrow\right| + \rho_{\downarrow\uparrow} \left|\downarrow\right\rangle \left\langle\uparrow\right|$$

$$\rightarrow \hat{P}_{\uparrow}\hat{\rho}\hat{P}_{\uparrow} + \hat{P}_{\downarrow}\hat{\rho}\hat{P}_{\downarrow} = \rho_{\uparrow\uparrow} |\uparrow\rangle\langle\uparrow| + \rho_{\downarrow\downarrow} |\downarrow\rangle\langle\downarrow|$$

Replace 1-shot non-selective measurement (decoherence) by dynamics!

40 > 40 > 42 > 42 > 2 > 900

Dynamical Non-Sel. Measurement, Decoherence

Time-continuous (dynamical) measurement of $\hat{A} = \sum_{k} A_{k} \hat{P}_{k}$:

$$d\hat{\rho}/dt = -\frac{1}{2}[\hat{A},[\hat{A},\hat{\rho}]]$$

Solution:

$$[\hat{A}, [\hat{A}, \hat{\rho}]] = \sum_{k} A_k^2 \hat{P}_k \hat{\rho} + \sum_{k} A_k^2 \hat{\rho} \hat{P}_k - 2 \sum_{k,l} A_k A_l \hat{P}_k \hat{\rho} \hat{P}_l$$

$$d(\hat{P}_n\hat{\rho}\hat{P}_m)/dt = -\frac{1}{2}\hat{P}_n[\hat{A},[\hat{A},\hat{\rho}]]\hat{P}_m = -\frac{1}{2}(A_m - A_n)^2(\hat{P}_n\hat{\rho}\hat{P}_m)$$

Off-diagonals \rightarrow 0, diagonals=const

Example:
$$\hat{A} = \hat{\sigma}_z$$
, $d\hat{\rho}/dt = -\frac{1}{2}[\hat{\sigma}_z, [\hat{\sigma}_z, \hat{\rho}]]$

$$\hat{\rho}(t) = \rho_{\uparrow\uparrow}(0) |\uparrow\rangle\langle\uparrow| + \rho_{\downarrow\downarrow}(0) |\downarrow\rangle\langle\downarrow| + e^{-2t}\rho_{\uparrow\downarrow}(0) |\uparrow\rangle\langle\downarrow| + e^{-2t}\rho_{\downarrow\uparrow}(0) |\downarrow\rangle\langle\uparrow| \rightarrow \rho_{\uparrow\uparrow}(0) |\uparrow\rangle\langle\uparrow| + \rho_{\downarrow\downarrow}(0) |\downarrow\rangle\langle\downarrow|$$

Master Equations

General non-unitary (but linear!) quantum dynamics:

$$d\hat{
ho}/dt = \mathcal{L}\hat{
ho}$$

Lindblad form — necessary and sufficient for consistency:

$$d\hat{
ho}/dt = -i[\hat{H},\hat{
ho}] + \left(\hat{L}\hat{
ho}\hat{L}^{\dagger} - \frac{1}{2}\hat{L}^{\dagger}\hat{L}\hat{
ho} - \frac{1}{2}\hat{
ho}\hat{L}^{\dagger}\hat{L}\right) + \dots$$

If
$$\hat{L} = \hat{L}^{\dagger} = \hat{A}$$
:

$$d\hat{\rho}/dt = -i[\hat{H},\hat{\rho}] - \frac{1}{2}[\hat{A},[\hat{A},\hat{\rho}]]$$

Decoherence (non-selectiv measurement) of \hat{A} competes with \hat{H} . General case $\hat{H} \neq 0, \hat{L} \neq \hat{L}^{\dagger}$: untitary, decohering, dissipative, pump mechanisms compete.



1-Shot Selective Measurement, Collapse

Measurement of
$$\hat{A} = \sum_n A_n \hat{P}_n$$
; $\sum_n \hat{P}_n = \hat{I}$, $\hat{P}_n \hat{P}_m = \delta_{nm} \hat{I}$ General (mixed state) and the special case (pure state), resp. mixed state: pure state, $\hat{P}_n = |n\rangle \langle n|$: $\hat{\rho} \to \frac{\hat{P}_n \hat{\rho} \hat{P}_n}{p_n} \equiv \hat{\rho}_n$ $|\psi\rangle \to |n\rangle \equiv |\psi_n\rangle$ with-prob. $p_n = \operatorname{tr}(\hat{P}_n \hat{\rho})$ with-prob. $p_n = |\langle n|\psi\rangle|^2$

Selective measurement is refinement of non-selective.

Mean of conditional states = Non-selective post-measurement state:

$$\begin{aligned} \mathbf{M}\hat{\rho}_{n} &= \sum_{n} p_{n} \hat{\rho}_{n} = \\ &= \sum_{n} \hat{P}_{n} \hat{\rho} \hat{P}_{n} &= \sum_{n} \hat{P}_{n} \left| \psi \right\rangle \left\langle \psi \right| \hat{P}_{n} \end{aligned}$$

Replace 1-shot selective measurement (collapse) by dynamics!

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Dynamical Non-selective Measurement, Collapse

Take pure state 1-shot measurement of $\hat{A} = \sum_n A_n |n\rangle \langle n|$ and expand it for asymptotic long times:

$$|\psi(0)\rangle$$
 evolves into $|\psi(t)\rangle \rightarrow |n\rangle$

Construct a (stationary) stochastic process $|\psi(t)\rangle$ for t>0 such that for any initial state $|\psi(0)\rangle$ the solution walks randomly into one of the orthogonal states $|n\rangle$ with probability $p_n=|\langle n|\psi(0)\rangle|^2!$ There are ∞ many such stochastic processes $|\psi(t)\rangle$. Luckily, for

$$\hat{
ho}(t) = \mathbf{M} \ket{\psi(t)} \bra{\psi(t)}$$

we have already constructed a possible non-selective dynamics, recall:

$$d\hat{\rho}/dt = -\frac{1}{2}[\hat{A},[\hat{A},\hat{\rho}]]$$

This is a major constraint for the process $|\psi(t)\rangle$. Infinite many choices still remain.

Dynamical Collapse: Diffusion or Jump

Consider the dynamical measurement of $\hat{A} = \sum_{n} A_n |n\rangle \langle n|$, described by dynamical decoherence (master) equation:

$$d\hat{
ho}/dt = -\frac{1}{2}[\hat{A}, [\hat{A}, \hat{
ho}]]$$

Construct stochastic process $|\psi(t)\rangle$ of dynamical collapse satisfying the master equation by $\hat{\rho}(t) = \mathbf{M} |\psi(t)\rangle \langle \psi(t)|$.

• Gisin's Diffusion Process (1984):

$$d\ket{\psi}/dt = -i\hat{H}\ket{\psi} - \frac{1}{2}(\hat{A} - \langle \hat{A} \rangle)^2\ket{\psi} + (\hat{A} - \langle \hat{A} \rangle)\ket{\psi}w_t$$

 w_t : standard white-noise; $\mathbf{M}w_t=0, \ \mathbf{M}w_tw_s=\delta(t-s)$

• Diosi's Jump Process (1985/86):

$$d\ket{\psi}/dt = -i\hat{H}\ket{\psi} - \frac{1}{2}(\hat{A} - \langle \hat{A} \rangle)^2\ket{\psi} + \frac{1}{2}\langle(\hat{A} - \langle \hat{A} \rangle)^2\rangle\ket{\psi}$$

jumps
$$|\psi(t)\rangle o {\sf const.} imes (\hat{\it A} - \langle \hat{\it A} \rangle) \, |\psi(t)\rangle$$
 at rate $\langle (\hat{\it A} - \langle \hat{\it A} \rangle)^2 \rangle$

Dynamical Collapse: Diffusion or Jump - Proof

• Gisin's Diffusion Process (1984):

$$d\ket{\psi}/dt = -i\hat{H}\ket{\psi} - \frac{1}{2}(\hat{A} - \langle\hat{A}\rangle)^2\ket{\psi} + (\hat{A} - \langle\hat{A}\rangle)\ket{\psi}w_t$$

 w_t : standard white-noise; $\mathbf{M}w_t=0, \ \mathbf{M}w_tw_s=\delta(t-s)$

• Diosi's Jump Process (1985/86):

$$d |\psi\rangle / dt = -i\hat{H} |\psi\rangle - \frac{1}{2}(\hat{A} - \langle \hat{A} \rangle)^2 |\psi\rangle + \frac{1}{2}\langle (\hat{A} - \langle \hat{A} \rangle)^2 \rangle |\psi\rangle$$

jumps
$$|\psi(t)\rangle \to \text{const.} \times (\hat{A} - \langle \hat{A} \rangle) |\psi(t)\rangle$$
 at rate $\langle (\hat{A} - \langle \hat{A} \rangle)^2 \rangle$

If $[\hat{H}, \hat{A}] = 0$, prove:

•
$$\hat{
ho}(t) = \mathbf{M} \ket{\psi(t)} \bra{\psi(t)}$$
 satisfies $d\hat{
ho}/dt = -\frac{1}{2} [\hat{A}, [\hat{A}, \hat{
ho}]]$

- $|\psi(t)\rangle \rightarrow |n\rangle$
- $|n\rangle$ occurs with $p_n = |\langle n|\psi(0)\rangle|^2$

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Revisit Early Motivations 1970-1980's

Interpretation of ψ is statistical.

Sudden 'one-shot' collapse $\psi \to \psi_n$ is central.

- If collapse takes time? Why not!
- Hunt for a math model (Pearle, Gisin, Diosi) Too many models!
- New physics?
 - No, it's standard physics of real time-continuous measurement (monitoring).
 - Yes, it's new!
 - to add universal non-unitary modifications to QM
 - to replace von Neumann statistical interpretation