

On canonical coupling of classical geometry to quantum matter

1. Introduction

$$R_{ab} - \frac{1}{2}g_{ab}R = \frac{8\pi G}{c^4}T_{ab}$$

1.1 classical T , classical R — fine.

1.2 quantum T , quantum R — nobody knows.

1.3 quantum T , classical R — try it canonically!

2. Poisson, Dirac, and Alexandrov brackets

3. Coarse-grained Alexandrov equation

4. Application to gravity

5. Fundamental uncertainty of geometry, decoherence

6. Conclusion

6.1 Semiclassical evolution must be coarse-grained

6.2 Classical part, too, gets noisy.

6.3 Reasonable decoherence obtained.

6.4 Covariant coarse-graining?

2.1 Classical: Poisson bracket, Liouville equation of motion.

$$\{A, B\}_P \equiv \sum_n \left(\frac{\partial A}{\partial q_n} \frac{\partial B}{\partial q_n} - \frac{\partial A}{\partial p_n} \frac{\partial B}{\partial p_n} \right)$$
$$\dot{\rho} = \{H, \rho\}_P$$

where $\rho(q, p) \geq 0$, $\int \rho(q, p) dq dp = 1$.

2.2 Quantum: Dirac bracket, Schrodinger (von Neumann) equation.

$$\{A, B\}_D \equiv -\frac{i}{\hbar} [A, B].$$
$$\dot{\rho} = \{H, \rho\}_D$$

where $\rho \geq 0$, $\text{tr} \rho = 1$.

2.3 Semiclassical: Alexandrov (1981) bracket and equation.

$$\{A, B\}_A = \{A, B\}_P + \frac{1}{2}\{A, B\}_D - \frac{1}{2}\{B, A\}_D$$
$$\dot{\rho} = \{H, \rho\}_A$$

where $\rho(q, p) \geq 0$, $\text{tr} \int \rho(q, p) dq dp = 1$.

3. Coarse-grained Alexandrov equation.

$$H = H(q_1, q_2, p_1, p_2) = H_1(q_1, p_1) + H_2(q_2, p_2) + H_I(q_1, q_2, p_1, p_2)$$

where (q_1, p_1) are quantum, (q_2, p_2) are classical.

$$H_I(q_1, p_1, q_2, p_2) = \sum_{\alpha} J_1^{\alpha}(q_1, p_1) J_2^{\alpha}(q_2, p_2), \quad (J_1^{\alpha}, J_2^{\alpha} \neq 1).$$

Noisy interaction:

$$H_I^{noise}(q_1, p_1, q_2, p_2) = \left(J_1^{\alpha}(q_1, p_1) + \delta J_1^{\alpha}(t) \right) \left(J_2^{\alpha}(q_2, p_2) + \delta J_2^{\alpha}(t) \right)$$

where $\delta J_1, \delta J_2$ are *classical* noises.

$$H^{noise} = H_1 + H_2 + H_I^{noise}$$

$$\dot{\rho} = \left\langle \{ H^{noise}, \rho \}_A \right\rangle_{noise}.$$

Choose white-noises:

$$\left\langle \delta J_n^{\alpha}(t') \delta J_n^{\beta}(t) \right\rangle_{noise} = \frac{1}{2} \Lambda_n^{\alpha\beta} \delta(t' - t), \quad n = 1, 2.$$

$$\begin{aligned} \dot{\rho} = & -\frac{i}{\hbar} [H, \rho] + \frac{1}{2} \{H, \rho\}_P - \frac{1}{2} \{\rho, H\}_P \\ & - \frac{1}{4} \sum_{\alpha, \beta} \Lambda_2^{\alpha\beta} [J_1^{\alpha}, [J_1^{\beta}, \rho]] + \frac{1}{4} \sum_{\alpha, \beta} \Lambda_1^{\alpha\beta} \{J_2^{\alpha}, \{J_2^{\beta}, \rho\}_P\}_P \end{aligned}$$

Lindblad's condition (1976) for positivity: $\Lambda_1 \Lambda_2 = I$.

4. Application to gravity.

$(q_1, p_1) \equiv (q, p)$ *quantized non – relativistic matter*

$(q_2, p_2) \equiv (\phi, \pi)$ *weak classical gravitational field*

$\phi \equiv \frac{1}{2}c^2(g_{00} - 1)$ *classical Newtonian potential*

$f = T_{00}/c^2$ *operator of mass distribution*

$$H(q, p, \phi, \pi) = H_m(q, p) + \frac{1}{8\pi G} \int_r \left(\frac{1}{c^2} \pi^2 + |\nabla \phi|^2 \right) + \int_r \phi(r) f(r)$$

Coarse-grained Alexandrov equation:

$$\begin{aligned} \dot{\rho} = & -\frac{i}{\hbar} [H_m, \rho] - \frac{i}{\hbar} \int_r \phi(r) [f(r), \rho] \\ & - \frac{1}{4\pi G} \int_r \left(\frac{1}{c^2} \pi(r) \frac{\delta \rho}{\delta \phi(r)} + \Delta \phi(r) \frac{\delta \rho}{\delta \pi(r)} \right) + \frac{1}{2} \int_r [f(r), \frac{\delta \rho}{\delta \pi(r)}]_+ \\ & - \frac{1}{4} \int_r \int_{r'} \lambda(r, r') [f(r), [f(r'), \rho]] + \frac{1}{4} \int_r \int_{r'} \lambda^{-1}(r, r') \frac{\delta^2 \rho}{\delta \pi(r) \delta \pi(r')}. \end{aligned}$$

Reduced dynamics of $\rho_m = \int \int \rho(\phi, \pi) \mathcal{D}\phi \mathcal{D}\pi$ in Newtonian approximation.
Drop term $1/c^2$ and assume ρ_m determines Newton potential ϕ :

$$\int \rho(\phi, \pi) \mathcal{D}\pi = \prod_r \delta \left(\phi(r) + \int_{r'} \frac{G/2}{|r - r'|} [f_+(r') + f_-(r')] \right) \rho_m.$$

Subscripts + and – assure ρ_m 's hermiticity.

Integrate both sides over ϕ, π !

$$\dot{\rho}_m = -\frac{i}{\hbar} [H_m + H_g, \rho_m] - \frac{1}{4} \int_r \int_{r'} \lambda(r, r') [f(r), [f(r'), \rho_m]]$$

$$H_g = -\frac{G}{2} \int_r \int_{r'} \frac{f(r) f(r')}{|r - r'|}.$$

5. Fundamental uncertainty of geometry, decoherence

$$\begin{aligned}\langle \delta f(r', t') \delta f(r, t) \rangle_{noise} &= \frac{1}{2} \lambda^{-1}(r', r) \delta(t' - t), \\ \langle \delta \phi(r', t') \delta \phi(r, t) \rangle_{noise} &= \frac{\hbar^2}{2} \lambda(r', r) \delta(t' - t).\end{aligned}$$

In Newtonian approximation: $\Delta \langle \phi(r) \rangle = 4\pi G \langle f(r) \rangle$.

$$\Delta \Delta' \lambda(r, r') = (4\pi G)^2 \lambda^{-1}(r, r')$$

Diósi and Lukács (1987):

$$\lambda(r, r') = (G/\hbar) |r - r'|^{-1}$$

Diósi (1987):

$$\dot{\rho}_m = -\frac{i}{\hbar} [H_m + H_g, \rho_m] - \frac{1}{4} \int_r \int_{r'} \frac{G/\hbar}{|r - r'|} [f(r), [f(r'), \rho_m]].$$

Ghirardi et al. (1990): cutoff $10^{-5} cm$.

Decoherence at characteristic time $\tau = \hbar/\Delta U_G$,

where ΔU_G : quantum spread of the Newtonian self-energy.

Pearle and Squires (1995): Verification in proton-decay experiments(!?)

References

- [1] I.V. Aleksandrov, Z.Naturforsch. **36A**, 902 (1981).
- [2] G.Lindblad, Commun.Math.Phys. **48**, 119 (1976).
- [3] L.Diósi and B.Lukács, Annln.Phys. **44**, 488 (1987).
- [4] L.Diósi, Phys.Lett. **120A**, 377 (1987); Phys.Rev. **A40**, 1165 (1989).
- [5] G.C.Ghirardi, R.Grassi and P.Pearle, Phys.Rev. **A42**, 1057 (1990).
- [6] E.Squires and P.Pearle, Phys.Rev.Lett. **73**, 1 (1994).